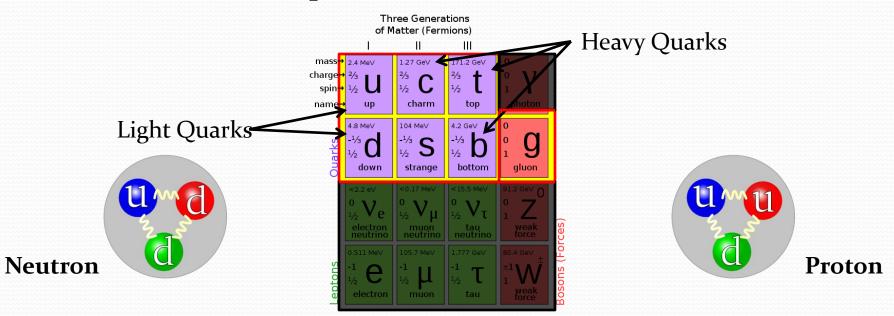


## Quarks and the Strong Force

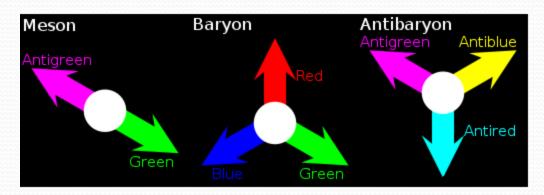
- The strong nuclear force governs the stability of hadrons.
- Predicted by theory, early pp collisions made first confirmation of quarks:



## Quantum Chromo-dynamics

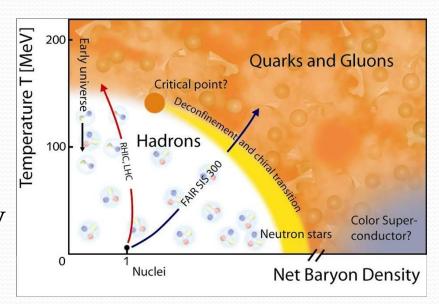
- The strong force is governed by Quantum Chromo-Dynamics (QCD), where interactions are mediated by gluons.
- Each quark and gluon has a color charge: red, green, blue.
- QCD gives two unusual properties of quarks and gluons:

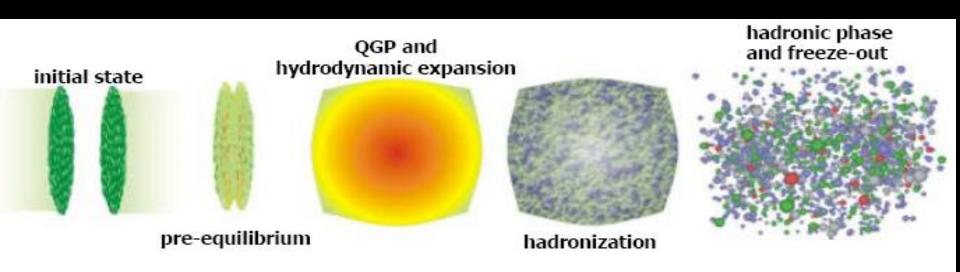
  - Asymptotic freedom



#### **Extreme Conditions**

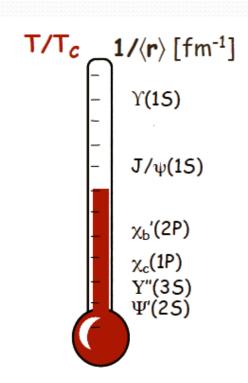
- QCD predicts the new deconfined state of matter formed around T<sub>c</sub> ≅ 180 MeV (~10<sup>12</sup> K): the Quark-Gluon Plasma (QGP).
- The creation and exploration of the QGP has been a primary science goal of relativistic heavy-ion colliders.
- However, the QGP is a transient state, existing for only ~10<sup>-22</sup> s. A probe is necessary





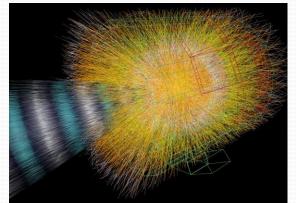
## Heavy Quarkonia

- Bound states of hidden flavor heavy quarkonia do not melt until well above  $T_c$ .
- These can be used as signatures of the QGP
  - Well studied cc̄ -> J/ψ
  - Less studied  $b\bar{b} \rightarrow \Upsilon$
- The complex interactions of the QGP requires a strong theoretical understanding in order to interpret experimental results.



## Heavy Quarkonia: Y

- The  $\Upsilon$  bound state is less studied than its cousin  $J/\psi$ .
- However, this particle provides a better signature of the QGP.
  - Survives up to 3-4 T<sub>C</sub>
  - Y yields unchanged during latter, cooler (hadronic) phases
  - True probe of the first  $\sim 10 \text{ fm/c} (\sim 10^{-23} \text{ s})$  of the QGP



### Y In-medium

- Primordial production of Y is affected by cold nuclear matter (CNM) effects: nuclear absorption and nuclear shadowing.
- Produced in N-N collisions at high energy, Υ undergoes several effects while in the QGP:
  - Suppression from dissociation due to scattering
  - Regeneration of Υ from recombined b b
  - Decay of higher bound states to Υ

Particle	Percent of Primordial Y(1s)
$X_b(1p)$ decay to $Y(1s)$	27%
$X_b(2p)$ decay to $\Upsilon(1s)$	10%
$\Upsilon(2s)$ [Y'] Decay to $\Upsilon(1s)$	11%
$\Upsilon(3s) [\Upsilon'']$ Decay to $\Upsilon(1s)$	1%

## **Binding Models**

 Heavy quarkonia resonances occur through one of two binding scenarios:

#### **Weak Binding**

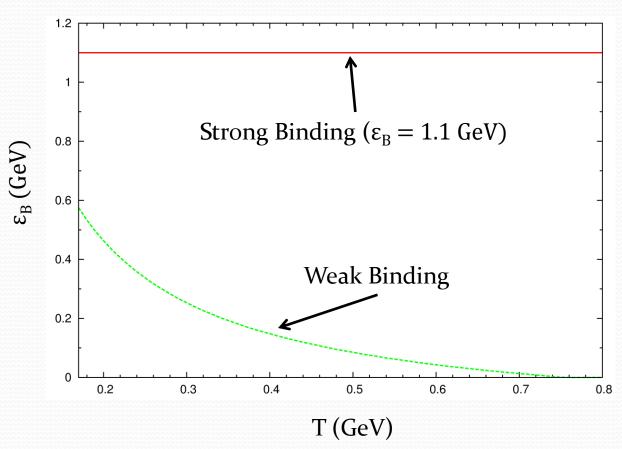
 In-medium bound state mass, binding energy, and unbound masses, satisfying:

$$\varepsilon_B(T) = 2m_b - m_Y$$

#### **Strong Binding**

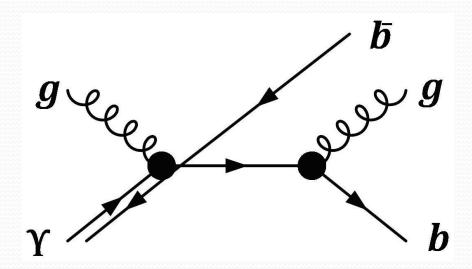
• Fixed, vacuum bound state mass with  $m_b = 5.280 \text{ GeV/c}^2$ .

## Y Binding Energy

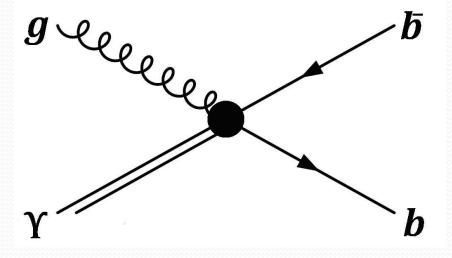


### Dissociation Mechanisms

 Each binding scenario has a corresponding dissociation mechanism that is most efficient:



Weak Binding – Quasifree Dissociation



Strong Binding – Gluo Dissociation

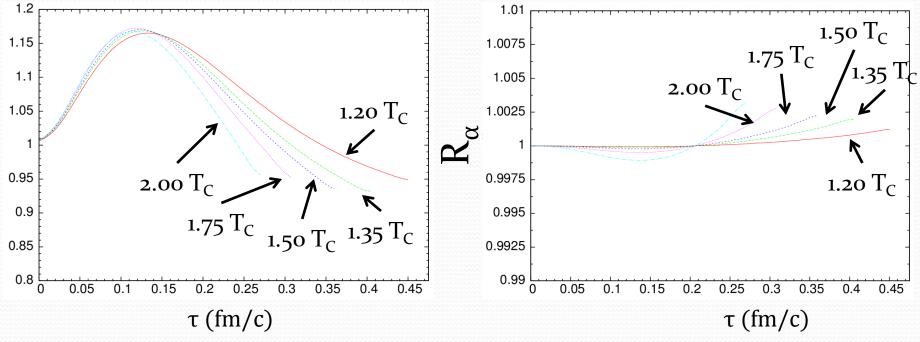
#### **Euclidean Correlator Ratios**

- Euclidean correlator ratios indicate in-medium effects on particle spectral functions:
- This property of Υ states is related to their spectral functions, readily calculable in numerical lattice QCD (lQCD) calculations.

$$R_{\alpha}(\tau,T) = \frac{G_{\alpha}(\tau,T)}{G_{\alpha}^{rec}(\tau,T)} = \frac{\int_{0}^{\infty} \sigma_{\alpha}(\omega,T) K(\omega,\tau,T) d\omega}{G_{\alpha}^{rec}(\tau,T)}$$

- The denominator is of the same form, evaluated at low temperature.
- Lattice QCD calculations indicate that  $R_{\alpha}$  remains close to unity.

## Correlator Ratios: New Insight



- The strong binding scenario remains close to unity (within 1%).
- This motivates necessity to revisit the strong binding scenario as the relevant model of  $\Upsilon$  in the QGP.

## Improvements and Updates

- In [L. Grandchamp et. al. PRC 73 (2006)], a study of Υ production at RHIC and LHC was made using these models.
- Better understanding of the physics of Y production and correlator ratio calculations have motivated an necessary update
- This update will yield a better insight into bottomonia in the QGP.

### Collision Model: Fireball

• A kinetic-theory rate-equation approach was used to model the collision and Y production

$$\frac{dN_{Y}}{d\tau} = -\Gamma_{Y} \left( N_{Y} - N_{Y}^{eq} \right)$$

- The fireball evolves as a function of time, allowing implementation of temperature dependence.
- CNM effects, suppression, regeneration, and feed-down included.
- For simplification, feed-down from  $\Upsilon$ " is ignored, and  $\chi_b(1p)$  and  $\chi_b(2p)$  states are not differentiated.

#### **Nuclear Modification Factor**

 The nuclear modification factor (R<sub>AA</sub>) is used to characterize particle production in NN collisions with respect to pp collisions.

$$R_{AA}^{Total} = rac{N_{
m Y}^{Total}}{N_{collision} rac{\sigma_{pp 
ightarrow 
m Y}}{\sigma_{pp}^{inelastic}}}$$

 This is a valuable tool used to analyze experimental data, and to compare theory to experiment.

## Probing the QGP in Experiment:

 The nuclear modification factor was calculated for two different cases:

# Reletavistic Heavy Ion Collisder (RHIC):

- AuAu collisions at  $E_{CMS}$ = 200 GeV.
- Used  $\sigma_{\text{Nuc. Abs.}} = [2.0, 3.1] \text{ mb}$
- Cannot (yet) resolve  $\Upsilon$ ,  $\Upsilon'$ , and  $\Upsilon''$  states separately.
- Give combined  $\Upsilon(1s+2s+3s)$

#### $R_{AA}$

Compare to STAR data

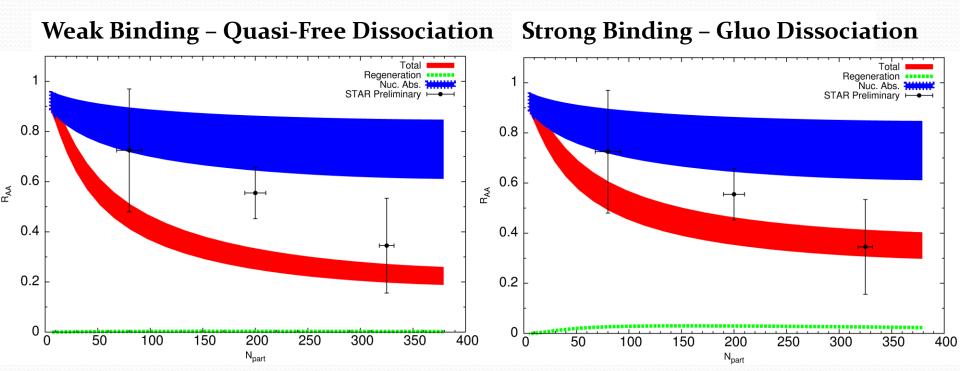
#### <u>Large Hadron Collider</u> (<u>LHC</u>):

- PbPb collisions at  $E_{CMS}$  = 2.76 TeV.
- Used  $\sigma_{\text{Nuc. Abs.}} = [0.0, 2.0] \text{ mb}$
- Give  $\Upsilon(1s)$  R<sub>AA</sub>
- Compare to CMS data



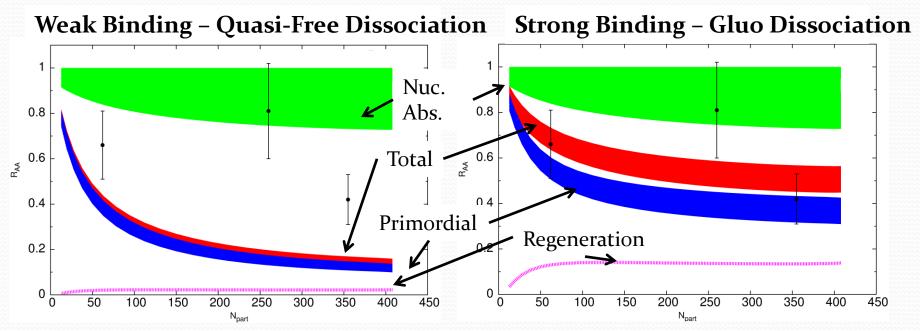
 $R_{AA}$  for higher Y states also calculated, though data only exists for these

### Results: RHIC



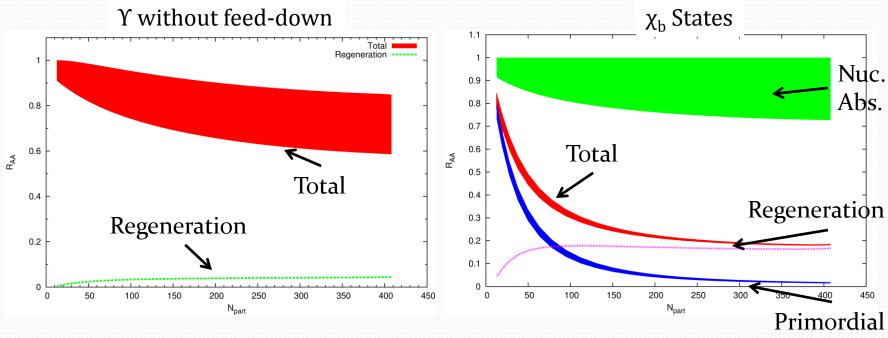
- The strong binding scenario gives better agreement with data
- The regeneration component is negligible for both binding scenarios

### Results: LHC



- The strong binding scenario shows much better agreement with data.
- Regeneration becomes significant at LHC
- Nuclear absorption is inconclusive

## Results: LHC (cont.)



- As shown, the regeneration component is significant in the higher states.
- The suppression of the  $\Upsilon(1s)$  state alone is 60-85% depending on nuclear absorption. Yet is ~45-55% including feed-down.
- Right shows large suppression of  $\chi_b$  state, with dominant regeneration component.

## Further Improvements

- The following would improve quality of checks between theory and experiment:
  - Increased statistics and data collection capabilities
  - Resolving individual Υ(nS) states at RHIC
  - Direct experimental measurement of Υ production crosssections at 2.76 TeV
  - Better theoretical calculations of  $\chi_h$  production cross sections
  - Improvement of Strong Binding model, implementing inmedium effects to the bound state mass. Especially important for higher bound states.
  - Applying models to calculations of p<sub>T</sub> spectra, adding another quantity to compare to data.

#### Conclusions:

- Modeled in medium Y particles with two binding scenarios: strong and weak
- Calculated Euclidean correlator ratios for each
- Used a kinetic-theory fireball model and rate-equation approach to calculate  $\Upsilon$   $R_{AA}$
- RHIC shows preference for strong binding model. Strong binding is appropriate model at LHC.
- Regeneration of Υ is significant at LHC
- More experimental data, and improved strong binding model will improve these.

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